## OSCILLATORY INTEGRALS OPEN PROBLEMS IN NUMBER THEORY SPRING 2018, TEL AVIV UNIVERSITY

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## 1. Oscillatory integrals

The Fourier transform  $\widehat{f}(k) = \int_{-\infty}^{\infty} f(x)e^{-2\pi ikx}dx$  is an example of an oscillatory integral. The general form that we take is

$$I(\lambda) := \int_{-\infty}^{\infty} A(x) e^{i\lambda\phi(x)} dx$$

where the phase function  $\phi(x)$  is assumed to be real, and the amplitude A(x) is assumed to be compactly supported, or at least in the Schwartz space S.

We will need to analyze a number of such oscillatory integrals and in particular understand their decay at infinity. The trivial bound is  $O(\int |A(x)|dx)$ , and we want to have some cancellation as  $\lambda \to \infty$ .

1.1. The principle of NON-stationary phase. Suppose first that A is smooth, and either

 $\begin{array}{l} 1) \ |\phi'| \geq 1 \\ \text{or} \end{array}$ 

2) if A is compactly supported we may simply assume that  $\phi' \neq 0$  has no critical points on the support of A.

Then as  $\lambda \to +\infty$ ,

$$I(\lambda) \ll_N \frac{1}{\lambda^N}, \quad \forall N \ge 1$$

(the implied constants depend on A and  $\phi$ ).

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*Proof.* The proof is by repeated integration by parts. Lets take the case that A is compactly supported. Define differential operators acting on functions supported in supp(A):

$$L = \frac{1}{i\phi'(x)}\frac{d}{dx}, \qquad L^T = -\frac{d}{dx} \circ \frac{1}{i\phi'(x)}$$

(this makes sense since  $\phi' \neq 0$  on the support of A). Integration by parts shows that for any  $f, g \in \mathcal{S}$ ,

$$\int_{-\infty}^{\infty} (Lf)(x)g(x)dx = \int_{-\infty}^{\infty} f(x)(L^{T}g)(x)dx$$

Moreover a computation shows that  $L(e^{i\lambda\phi(x)}) = \lambda e^{i\lambda\phi(x)}$ , so that

$$\mathcal{L}^{N}(e^{i\lambda\phi(x)}) = \lambda^{N} e^{i\lambda\phi(x)}, \qquad \forall N \ge 1.$$

Now we have

$$I(\lambda) = \int_{-\infty}^{\infty} e^{i\lambda\phi(x)} A(x) dx = \int_{-\infty}^{\infty} \frac{1}{\lambda^N} L^N(e^{i\lambda\phi(x)}) A(x) dx = \frac{1}{\lambda^N} \int_{-\infty}^{\infty} e^{i\lambda\phi(x)} (L^T)^N(A)(x) dx$$

and taking absolute values gives

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$$|I(\lambda)| \le \frac{1}{\lambda^N} \int_{-\infty}^{\infty} |(L^T)^N(A)(x)| dx = \frac{C_N}{\lambda^N}$$

with  $C_N = \int_{-\infty}^{\infty} |(L^T)^N(A)(x)| dx < \infty$  since  $(L^T)^N(A)$  is also smooth and compactly supported.

1.2. Van der Corput's Lemma. In the above, it was crucial to have A smooth in order to integrate by parts. When this does not hold, the result may fail. We will need the case  $A = \mathbf{1}_{[a,b]}$ , when

$$I(\lambda) = \int_{a}^{b} e^{i\lambda\phi(x)} dx$$

We first assume that  $\phi$  has no stationary (critical) points on [a, b]:

**Proposition 1.1** (Van der Corput's Lemma 1). Assume that  $\phi$  is smooth, that  $\phi' \neq 0$  on [a, b], and that  $\phi'$  is monotonic. Then

$$\Big|\int_a^b e^{i\lambda\phi(x)}dx\Big| \leq \frac{4}{\min_{x\in[a,b]}|\phi'(x)|}\cdot \frac{1}{|\lambda|}$$

*Proof.* Since  $\phi'$  is continuous and nonzero on [a, b], we have

$$c := \min_{x \in [a,b]} |\phi'(x)| > 0$$

Integrating by parts, we have

$$\begin{split} I(\lambda) &= \int_{a}^{b} e^{i\lambda\phi(x)} dx = \int_{a}^{b} e^{i\lambda\phi(x)} i\lambda\phi'(x) \cdot \frac{1}{i\lambda\phi'(x)} dx \\ &= e^{i\lambda\phi(x)} \frac{1}{i\lambda\phi'(x))} \Big|_{a}^{b} - \int_{a}^{b} e^{i\lambda\phi(x)} \frac{d}{dx} \Big(\frac{1}{i\lambda\phi'(x)}\Big) dx \end{split}$$

and therefore (WLOG take  $\lambda > 0$ )

$$|I(\lambda)| \le \frac{1}{\lambda} \Big| \frac{e^{i\lambda\phi(b)}}{\phi'(b)} - \frac{e^{i\lambda\phi(a)}}{\phi'(a)} \Big| + \frac{1}{\lambda} \int_a^b \Big| \frac{d}{dx} \Big\{ \frac{1}{\phi'(x)} \Big\} \Big| dx.$$

Since  $|\phi'| \ge c$  on [a, b], we have

$$\Big|\frac{e^{i\lambda\phi(b)}}{\phi'(b)}-\frac{e^{i\lambda\phi(a)}}{\phi'(a)}\Big|\leq \frac{1}{|\phi'(b)|}+\frac{1}{|\phi'(a)|}\leq \frac{2}{c}.$$

For the integral, since  $\phi'$  is monotonic (and nonzero), we have  $1/\phi'$  monotonic, so that  $\frac{d}{dx} \left\{ \frac{1}{\phi'(x)} \right\}$  has a fixed sign. Therefore

$$\int_{a}^{b} \left| \frac{d}{dx} \left\{ \frac{1}{\phi'(x)} \right\} \right| dx = \left| \int_{a}^{b} \frac{d}{dx} \left\{ \frac{1}{\phi'(x)} \right\} dx \right| = \left| \frac{1}{\phi'(b)} - \frac{1}{\phi'(a)} \right| \le \frac{2}{c}$$

as before. Thus we find

$$|I(\lambda)| \le \frac{4/c}{\lambda}.$$

We now turn to study the case when the phase function  $\phi$  has critical (stationary) points. Assume that all critical points are non-degenerate ( $\phi''(x) \neq 0$  if  $\phi'(x) = 0$ ). By subdividing the interval<sup>1</sup>, we can assume that either case 1 holds (no critical points) or  $\phi'' \neq 0$  on the entire interval [a, b]. In that case we have

**Proposition 1.2** (Van der Corput's Lemma 2). Let  $\phi$  be real valued and smooth on [a, b], with  $\phi'' \neq 0$  on [a, b]. Then

$$\left|\int_{a}^{b} e^{i\lambda\phi(x)} dx\right| \leq \frac{8}{\sqrt{\min_{x\in[a,b]} |\phi''(x)|}} \cdot \frac{1}{\sqrt{|\lambda|}}.$$

*Proof.* By replacing  $\phi$  by  $\phi/\min_{x\in[a,b]} |\phi''(x)|$  and  $\lambda$  by  $\min_{x\in[a,b]} |\phi''(x)| \cdot \lambda$ , we may assume that  $|\phi''(x)| \ge 1$ , and WLOG take  $\phi'' \ge 1$ .

Let  $c \in [a, b]$  be a point where the first derivative  $|\phi'|$  is minimal:  $|\phi'(c)| \leq |\phi'(x)|$  for all  $x \in [a, b]$ . Since the second derivative is non-vanishing, it cannot be the case that c is an interior local minimum/maximum of  $\phi'$ , and hence either  $\phi'(c) = 0$  or c is one of the endpoints a, b.

Assume first that  $\phi'(c) = 0$ , as in Figure 1. Then outside the interval  $(c-\delta, c+\delta)$ ,



FIGURE 1. A local minimum of  $|\phi'|$  where  $\phi'(c) = 0$ .

<sup>&</sup>lt;sup>1</sup>We assume that  $\phi$  has only finitely many critical points in [a, b], which would be the case if it was real analytic.

we have  $|\phi'(x)| \ge \delta$ , because we assume  $\phi''(x) \ge 1$  for all  $x \in [a, b]$ . Indeed, if  $c + \delta \leq x \leq b$  then

$$\phi'(x) = \phi'(x) - \phi'(c) = \int_{c}^{x} \phi''(t)dt \ge \int_{c}^{x} 1dt = x - c \ge \delta$$

with a similar argument if  $a \leq x \leq c - \delta$ . Now divide the range of integration [a, b]into three subintervals

$$\int_{a}^{b} e^{i\lambda\phi(x)} dx = \int_{a}^{c-\delta} + \int_{c-\delta}^{c+\delta} + \int_{c+\delta}^{b}$$

On the interval  $[a, c-\delta]$ , use  $|\phi'| \ge \delta$  and the fact that  $\phi'$  is monotonic increasing (since  $\phi'' \ge 1 > 0$ ) to invoke Proposition 1.1 giving

$$\int_{a}^{c-\delta} e^{i\lambda\phi(x)} dx \Big| \le \frac{4}{|\min(|\phi'(x)| : x \in [a, c-\delta])} \cdot \frac{1}{\lambda} \le \frac{4}{\delta} \cdot \frac{1}{\lambda}$$

with the same bound holding for  $\int_{c+\delta}^{b}$ .

On the middle interval  $(c - \delta, c + \delta)$ , just estimate trivially  $|e^{i\lambda\phi(x)}| \leq 1$  giving

$$\Big|\int_{c-\delta}^{c+\delta} e^{i\lambda\phi(x)}dx\Big| \le 2\delta$$

Thus we find

$$\Big|\int_{a}^{b} e^{i\lambda\phi(x)}dx\Big| \leq \frac{8}{\delta\lambda} + 2\delta.$$

Taking  $\delta = 2/\sqrt{\lambda}$  gives

$$\Big|\int_{a}^{b} e^{i\lambda\phi(x)} dx\Big| \le \frac{8}{\sqrt{\lambda}}.$$

It remains to deal with the case that c is one of the endpoints, say c = a and  $\phi'(a) \neq 0$ , say  $\phi'(x) \geq \phi'(a) > 0$  for all  $x \in [a, b]$ . Then as before  $\phi'(x) \geq \delta$  for  $x \in [a + \delta, b]$  since  $\phi'' \ge 1$ , and then the previous argument shows that

$$\left| \int_{a}^{b} e^{i\lambda\phi(x)} dx \right| \le \left| \int_{a}^{a+\delta} 1 dx \right| + \left| \int_{a+\delta}^{b} e^{i\lambda\phi(x)} dx \right| \le \delta + \frac{4}{\delta\lambda} \le \frac{4}{\sqrt{\lambda}}$$

$$ag \ \delta = 2/\sqrt{\lambda}. \qquad \Box$$

on taking  $\delta = 2/\sqrt{\lambda}$ .

Remark: A similar result holds for the case of degenerate critical points, if we assume that for some  $k \ge 2$ , we have  $\phi^{(k)} \ne 0$  on [a, b]:

**Proposition 1.3.** There is an absolute constant  $c_k > 0$  so that for all smooth, real valued  $\phi$  with  $\phi^{(k)} \neq 0$  on [a, b],

$$\left| \int_{a}^{b} e^{i\lambda\phi(x)} dx \right| \leq \frac{c_{k}}{(\min_{x\in[a,b]} |\phi^{(k)}(x)|)^{1/k}} \cdot \frac{1}{|\lambda|^{1/k}}.$$

**Exercise 1.** The Bessel function  $J_0(z) = \sum_{m=0}^{\infty} \frac{(-1)^m}{(m!)^2} \left(\frac{z}{2}\right)^{2m}$  admits an integral representation

$$J_0(z) = \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{-iz\sin t} dt$$

Show that as  $z \to +\infty$ ,  $J_0(z) \ll 1/\sqrt{z}$ .

We now include an amplitude

**Corollary 1.4.** Assume that  $A(x) \in C^1[a, b]$  is differentiable, that  $\phi$  is smooth and real valued, and that  $\phi^{(k)} \neq 0$  on [a, b] (and if k = 1 also that  $\phi'$  is monotonic). Then

$$\left|\int_{a}^{b} A(x)e^{i\lambda\phi(x)}dx\right| \leq \frac{c_{k}}{(\min_{x\in[a,b]}|\phi^{(k)}(x)|)^{1/k}} \cdot \left(|A(b)| + \int_{a}^{b} |A'(t)|dt\right) \frac{1}{|\lambda|^{1/k}}.$$

*Proof.* For notational simplicity we treat the case k = 2. Denote

$$J_{\lambda}(t) = \int_{a}^{t} e^{i\lambda\phi(x)} dx$$

So that  $J_{\lambda}(t) = e^{i\lambda\phi(t)}$ . Integrating by parts, we have

$$\int_{a}^{b} A(x)e^{i\lambda\phi(x)}dx = A(t)J_{\lambda}(t)\Big|_{a}^{b} - \int_{a}^{b} A'(t)J_{\lambda}(t)dt$$

Using our results for  $A \equiv 1$  for  $J_{\lambda}(t)$ , we have

$$\left|A(t)J_{\lambda}(t)\right|_{a}^{b} = \left|A(b)\int_{a}^{b} e^{i\lambda\phi(x)}dx\right| \le |A(b)|\frac{8}{\sqrt{\min_{x\in[a,b]}|\phi''(x)|}} \cdot \frac{1}{\sqrt{|\lambda|}}$$

and

$$\int_{a}^{b} |A'(t)| |J_{\lambda}(t)| dt \leq \int_{a}^{b} |A'(t)| \frac{8}{\sqrt{\min_{x \in [a,b]} |\phi''(x)|}} \cdot \frac{1}{\sqrt{|\lambda|}} dt$$

**Corollary 1.5.** Assume that the amplitude A and the phase function  $\phi$  are smooth, and that  $\phi$  has finitely many critical points, all of them non-degenerate. Then

$$I(\lambda) \ll \frac{1}{\sqrt{\lambda}}, \quad \lambda \to +\infty$$

the implied constant depending on A and  $\phi$ .

Proof. To use van der Corput's Lemma, we use a smooth partition of unity to write

$$\mathbf{1}_{[0,2\pi]} = \sum_{j} \psi_{j}$$

where  $\psi_j$  are smooth, the support of each contains at most one of the critical points, and when the support does contain a critical point  $x_0$  (at which  $\phi''(x_0) \neq 0$ ), we take the support sufficiently small so that  $\phi'' \neq 0$  on all of  $\sup \psi_j$ , while the remaining  $\psi_j$  are supported away from the critical points. Hence we can write

$$I(\lambda) = \sum_{j} I_{j}(\lambda), \qquad I_{j}(\lambda) = \int_{0}^{2\pi} \psi_{j}(t) A(t) e^{i\lambda\phi(t)} dt.$$

To bound the integrals  $I_j(\gamma)$  where the support of  $\psi_j$  does not include any critical points, we use the principle of NON-stationary phase, with a smooth amplitude and a phase function so that  $|\phi'| \ge c_j > 0$  where  $c_j = \min(|\phi'(x)| : x \in \operatorname{supp} \psi_j)$  to bound

$$I_j(\lambda) \ll \frac{1}{|\xi|^N}, \quad \forall N \ge 1.$$

For j such that  $\operatorname{supp} \psi_j$  contains a critical point (unique by assumption), we can Corollary 1.4 with k = 2, since we have taken the support of such  $\psi_j$  so that  $\phi'' \neq 0$ on  $\operatorname{supp} \psi_j$ . Hence for such j we have

$$I_j(\lambda) \ll \frac{1}{\sqrt{\lambda}}.$$

Altogether we have proven  $I(\lambda) \ll \lambda^{-1/2}$ .

1.3. An asymptotic expansion: The method of stationary phase. It is possible to go from upper bounds to asymptotic expansions. I will quote a typical result (we will not explicitly use this).

**Theorem 1.6.** Assume that the amplitude  $A \in C_c^{\infty}(\mathbb{R})$  is smooth and compactly supported, and that the phase function (real valued and smooth) has a single critical point  $x_0 \in \text{supp } A$ , which is non-degenerate:  $\phi'(x_0) = 0$ ,  $\phi''(x_0) \neq 0$ . Then

$$I(\lambda) \sim e^{i\frac{\pi}{4}\operatorname{sign}(\phi''(x_0))} A(x_0) \sqrt{\frac{2\pi}{|\phi''(x_0)|}} \cdot \frac{e^{i\lambda\phi(x_0)}}{\sqrt{\lambda}}, \quad \text{as } \lambda \to +\infty.$$

1.4. An application: The Fourier transform of the unit disk. We want to bound the Fourier transform of the unit disk  $B(0,1) \subset \mathbb{R}^2$ : If  $\chi(x) = 1$ ,  $|x| \leq 1$ , and is zero otherwise, the Fourier transform is (we have dropped the factor of  $-2\pi$  from the exponent)

$$\widehat{\chi}(\xi) = \int_{|x| \le 1} e^{i\langle \xi, x \rangle} dx, \qquad \xi \in \mathbb{R}^2.$$

Proposition 1.7.

$$\widehat{\chi}(\xi) \ll \frac{1}{1+|\xi|^{3/2}}.$$

*Proof.* We convert the 2-dimensional integral to a one-dimensional integral, to which we can apply the van der Corput bound, by using Green's theorem. Recall that Green's theorem says that for a bounded planar domain D, with a piecewise smooth boundary  $\partial D$ , we have for  $A, B \in C^2(\mathbb{R}^2)$ ,

$$\int_{D} \left( \frac{\partial B}{\partial x} - \frac{\partial A}{\partial y} \right) dx \wedge dy = \oint_{\partial D} A dx + B dy$$

where the line integral over the boundary  $\partial D$  is taken counterclockwise.

For us, D = B(0,1) is the unit disk, with boundary  $\partial D = S^1$  the unit circle, and we want to find A, B so that

$$\frac{\partial B}{\partial x} - \frac{\partial A}{\partial y} = e^{i(ax+by)}, \quad \xi = (a,b) \neq 0.$$

A solution is to take

$$A=\frac{ibe^{i(ax+by)}}{|\xi|^2}, \quad B=\frac{-iae^{i(ax+by)}}{|\xi|^2}$$

Hence we find

$$\widehat{\chi}(\xi) = \frac{i}{|\xi|^2} \int_{\partial D} e^{i(ax+by)} \ (bdx - ady).$$

Parameterizing  $\partial D = S^1$  by arc-length:  $\gamma(t) = (\cos t, \sin t)$ , which runs counterclockwise if t runs from 0 to  $2\pi$ , we obtain

$$\widehat{\chi}(\xi) = \frac{i}{|\xi|^2} \int_0^{2\pi} e^{i(a\cos t + b\sin t)} (-b\sin t - a\cos t) dt = \frac{i}{|\xi|} I(|\xi|)$$

where

$$I(\lambda) = \int_0^{2\pi} A_{\xi}(t) e^{i\lambda\phi_{\xi}(t)} dt$$

is an oscillatory integral, with amplitude

$$A_{\xi}(t) = \langle \dot{\gamma}(t), \frac{\xi^{\perp}}{|\xi|} \rangle = \frac{-a\cos t - b\sin t}{\sqrt{a^2 + b^2}}$$

where  $\xi^{\perp} = (b, -a)$ , and phase function

$$\phi_{\xi}(t) = \langle \gamma(t), \frac{\xi}{|\xi|} \rangle = \frac{a \cos t + b \sin t}{\sqrt{a^2 + b^2}}.$$

We want to apply our estimates on oscillatory integrals to this case. Note that both the amplitude and the phase function depend on  $\xi/|\xi|$ , so it is important to make sure that our estimates our uniform in  $\xi/|\xi| \in S^1$ .

The phase function has two critical points  $\phi'(t) = \langle \dot{\gamma}(t), \xi/|\xi| \rangle = 0$ , when  $\dot{\gamma}(t) \perp \xi$ , i.e. when  $\dot{\gamma}(t) = \pm \xi^{\perp}/|\xi|$  (note  $\dot{\gamma}$  is a unit vector), with  $\xi^{\perp} = (b, -a)$ , say at  $t_0$  and therefore also at  $t_0 + \pi$ .

We claim that these critical points are non-degenerate: The second derivative of the phase function is

$$\phi''(t_0) = \langle \ddot{\gamma}(t_0), \xi \rangle.$$

Now further note that  $\langle \ddot{\gamma}, \dot{\gamma} \rangle = 0$  (which follows by direct computation or better yet by differentiating  $|\dot{\gamma}|^2 = 1$ , which is arc-length parameterization condition), and since also  $\xi \perp \dot{\gamma}(t_0)$ , we must have

$$\ddot{\gamma}(t_0) = \pm \kappa \frac{\xi}{|\xi|}$$

where  $\kappa = |\ddot{\gamma}(t_0)|$ ; note that here  $|\ddot{\gamma}| \equiv 1$  so  $\kappa = 1$  and hence<sup>2</sup>

$$\phi''(t_0) = \langle \ddot{\gamma}(t_0), \frac{\xi}{|\xi|} \rangle = \pm \langle \frac{\xi}{|\xi|}, \frac{\xi}{|\xi|} \rangle = \pm 1 \neq 0.$$

To use van der Corput's Lemma, we use a smooth partition of unity to write

$$\mathbf{1}_{[0,2\pi]} = \psi_1 + \psi_2 + \psi_3$$

where  $\psi_j$  are smooth,  $\psi_1$  is supported in say  $(t_0 - 0.1, t_0 + 0.1)$ ,  $\psi_1$  is supported in  $(t_0 + \pi - 0.1, t_0 + \pi + 0.1)$  and  $\psi_3$  is supported away from the critical points  $t_0, t_0 + \pi$ . This gives

with

$$I(\lambda) = I_1 + I_2 + I_3$$
$$I_j(\lambda) = \int_0^{2\pi} \psi_j(t) A(t) e^{i\lambda\phi(t)} dt.$$

<sup>&</sup>lt;sup>2</sup>For general D,  $|\ddot{\gamma}| = \kappa$  will be the curvature of  $\partial D$ ; if we assume that the boundary  $\partial D$  has nowhere vanishing curvature, then this computation shows that the critical points are all non-degenerate.

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For the integrals  $I_1$ ,  $I_2$  we can use van der Corput's Lemma to deduce that

$$|I_1(|\xi|)| \le \frac{8}{(\min_{x \in [t_0 - 0.1, t_0 + 0.1]} |\phi''(x)|)^{1/2}} \cdot \left(|A(t_0 + 0.1)| + \int_0^{2\pi} |A'(t)| dt\right) \frac{1}{|\xi|^{1/2}}.$$
  
We have

$$|A(t)| = |\langle \dot{\gamma}(t), \frac{\xi}{|\xi|} \rangle| \le |\dot{\gamma}| \cdot |\frac{\xi}{|\xi|} = 1$$

and likewise  $|A'(t)| \leq 1$ . Since  $\phi''(t_0) = \pm 1$ , we have

$$\min(|\phi''(t)| : t_0 - 0.1 < t < t_0 + 0.1) > c_0$$

for some positive constant  $c_0$ . Hence we deduce from van der Corput's Lemma that there is some C > 0, independent of  $\xi$ , so that

$$|I_1(|\xi|)|, |I_2(|\xi|)| \le \frac{C}{|\xi|^{1/2}}.$$

To bound the third integral  $I_3(\gamma)$ , we use the principle of NON-stationary phase, with a smooth amplitude and a phase function so that  $|\phi'| \ge c_3 > 0$  (uniformly in  $\xi$ ) to bound

$$I_3 \ll \frac{1}{|\xi|^N}, \quad \forall N \ge 1.$$

Altogether we obtain that uniformly in  $|\xi| \ge 1$ ,

$$I(\lambda) \le \frac{c_4}{\sqrt{\lambda}}$$
$$|\widehat{\chi}(\xi)| \le \frac{c_4}{|\xi|^{3/2}}.$$

and hence

$$\int_{|x|\leq 1}e^{i\langle x,\xi\rangle}d^3x\ll \frac{1}{1+|\xi|^2},\quad \xi\in\mathbb{R}^3.$$

Hint: Here we can directly evaluate the Fourier transform in elementary terms!